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## Study of Nonlinear Optical Properties of Disperse Yellow – 7 Dye-doped Polymer Films using CW Laser

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### ABSTRACT

The materials with an ability of self-focusing/defocusing through their light intensity-dependent refractive index and phase shift properties are potential candidates for all-optical photonic devices. Several organic materials including dye-doped polymer films are attracting many researchers due to their advantages in terms of enhanced efficiency, and acceptable mechanical properties to fabricate useful devices. In this paper, we have reported the nonlinear optical properties like nonlinear absorption and nonlinear refraction of an organic dye Disperse Yellow - 7 (DY-7) doped in Polymethyl methacrylate-methacrylic acid (PMMA-MA) polymer matrix using open aperture and closed aperture Z-scan experimental methods by means of low power continuous wave (CW) laser beam. The optical limiting properties of these films are also studied using Type 1 and Type 2 configurations at different input powers using continuous wave (CW) laser beams of 532 nm wavelength. The nonlinear refractive index  $n_2$ , nonlinear absorption coefficient  $\beta$ , changes in refractive index with input intensity, and the magnitude of third-order optical nonlinearity of the dye-doped film are experimentally determined. The sample is tested for optical power limiting in which the input limiting threshold and saturated output power level for both type 1 and type 2 optical limiting configurations are recorded. DY-7 dye-doped in PMMA-MA polymer film has shown saturation absorption at lower input irradiance and reverse saturation absorption at higher input irradiance along with thermal nonlinearity and hence found to be a potential candidate for the third harmonic property based photonic devices.

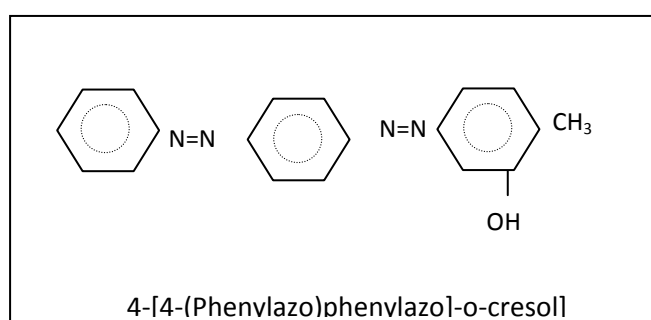
**Keywords:** Organic Nonlinear material, Dye-doped polymer films, Optical limiting, Open and closed aperture Z scan. DY-7 doped PMMA-MA films.

### 1. INTRODUCTION

The organic materials with an ability of self-focusing/defocusing through their light intensity-dependent refractive index and phase shift properties are potential materials in all-optical photonic devices have become an interesting topic of continuous research. Several organic materials including dye-doped polymer films are attracting many researchers due to their advantages in terms of enhanced efficiency, and acceptable mechanical properties to fabricate useful devices [1-5]. Many research results have been published on different organic nonlinear optical materials which include molecular single crystals, organic molecules in solution forms, and organic dyes and organometallics doped in different solid matrices. The mechanisms like Reverse Saturation Absorption (RSA), Two-Photon Absorption (TPA), nonlinear refraction, nonlinear scattering, and thermal nonlinearity are used in organic materials for optical power limiting. Similarly, the optical phase conjugation through four-

wave mixing is demonstrated in some of the gases, organic and inorganic crystals, and films using both pulsed and continuous wave (CW) lasers. Polymers, glasses and other types of solid matrices doped by nonlinear organic dyes are considered as promising materials for optical phase conjugation (OPC) due to their enhanced third-order nonlinear susceptibilities [6-10]. In this paper, we have reported the nonlinear optical properties like nonlinear absorption and nonlinear refraction of an organic dye Disperse yellow-7 (DY-7) doped in Polymethyl methacrylate-methacrylic acid (PMMA-MA) polymer matrix using open aperture and closed aperture Z-scan experimental methods by means of low power continuous wave (CW) laser beam. The optical limiting properties of these films are also studied using Type 1 and Type 2 configurations at different input powers using continuous wave (CW) laser beams of 532 nm wavelength. The nonlinear refractive index  $n_2$ , nonlinear absorption coefficient  $\beta$ , changes in refractive index with input intensity, and the magnitude of third-order optical nonlinearity of the dye-doped film are experimentally determined. In this paper, we have reported the nonlinear optical properties like nonlinear absorption and nonlinear refraction of an organic dye Disperse Yellow-7 (DY-7) doped in Polymethyl methacrylate-methacrylic acid (PMMA-MA) polymer matrix using open aperture and closed aperture Z-scan experimental methods by means of low power continuous wave (CW) laser beam. The optical limiting properties of these films are also studied using Type 1 and Type 2 configurations at different input powers using continuous wave (CW) laser beams of 532 nm wavelength. The nonlinear refractive index  $n_2$ , nonlinear absorption coefficient  $\beta$ , changes in refractive index with input intensity, and the magnitude of third-order optical nonlinearity of the dye-doped film are experimentally determined.

## 2. MOLECULAR STRUCTURE AND LINEAR OPTICAL PROPERTIES OF DY-7



**Figure 1:** Molecular structure of Disperse Yellow – 7.

Commercially available Disperse yellow-7 is purchased by Aldrich Chemical Company, USA. and further filtered and recrystallized using spectrograde ethanol. The research grade chloroform is used as the solvent. The thin films of DY-7 doped in polymer matrix by dissolving both DY-7 and PMMA-MA in specified quantity in chloroform. The solution mixtures are poured between two micro glass slides in a clean laboratory environment and thin films of both the dye samples of specified thickness are prepared using hot press technique [11]. Thin films of 10  $\mu\text{m}$  thickness of both DY-7 dye-doped polymer films are prepared between two glass slides separately. The thickness of the films is determined by gravimetric weighing method. The film thickness is determined to be 10  $\mu\text{m}$  for the solution molarity of 1 mM, 2 mM, & 5 mM. This is consequently verified by the cross-sectional studies of the film using Scanning Electron Microscope by mounting the film vertically to measure the thickness directly [12].

The linear absorption coefficient  $\alpha_0$  is determined for two wavelengths 1.06  $\mu\text{m}$  and 532 nm by using formula  $\alpha_0 = -\frac{1}{t} \ln \left[ \frac{1}{T} \right]$  ----- (1)

where (t) is the thickness of sample and T is the transmittance.

**Table 1:** Linear absorption coefficient and refractive index versus wavelength of DY-7.

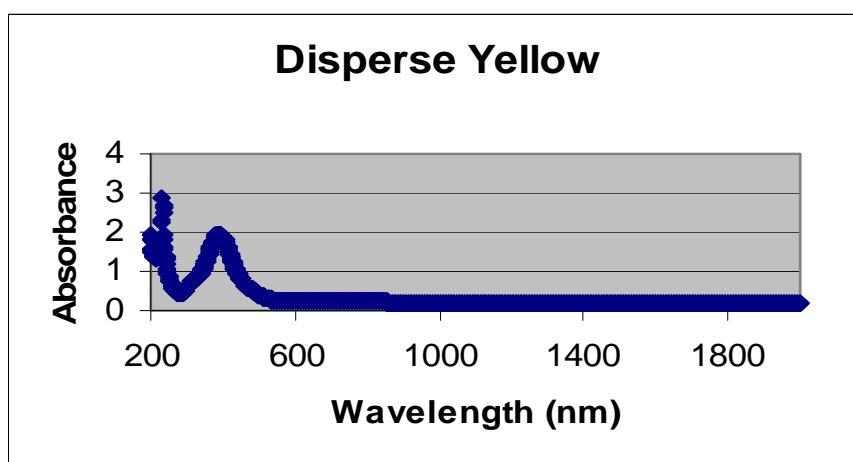
Wavelength (nm)	Thickness ( $\mu\text{m}$ )	Transmittance	$\alpha_0$ ( $\mu\text{m}$ ) <sup>-1</sup>	$n_0$
532	10	0.58	-0.05447226	3.12864
633	10	0.9	-0.010536	1.59543
1064	10	0.98	-0.00202026	1.22347

The refractive index  $n_0$  can be found from transmittance spectrum of the film according to the following equation

$$n_0 = \frac{1}{T} + \left[ \left( \frac{1}{T^2} - 1 \right) \right]^{1/2} \quad \text{----- (2)}$$

The linear absorption coefficient and refractive index are shown in Table 1.

The absorption spectra of the thin DY-7 dye-doped films is shown in figure 3 infer that transmission through these films is relatively low at below band gap region, indicating a high concentration of defects, free carriers. The transmittance increases abruptly in the short wavelengths which are due to the band edge absorption. A sudden increase at a particular wavelength indicates the presence of optical band gap in these samples. The incoming photons get sufficient energy to excite electrons from the valence band to the conduction band, which results in strong absorption in dye samples. The optical absorption edge is determined by the optical absorption, a simple and common method used that provides an explanation for the features concerning the band structure of the films [13-15].



**Figure 2 :** Ground state absorption spectrum of Disperse Yellow-7.

**Table 2 :** Physical parameters and Linear Absorption properties of DY-7 dye-doped film.

S. No.	Parameter/Property	DY-7 in PMMA-MA
1	Linear absorption range (bandwidth)	80 nm
2	Linear absorption peak (nm)	468 nm
3	Linear Transmission Range (nm)	528 – 1,600
4	Linear Transmittance (T) at 532 nm	0.58
5	Linear absorption coefficient ( $\alpha_0$ ) at 532 nm	-0.05447226
6	Linear Refractive index ( $n_0$ ) at 532 nm	3.12864
7	Dye concentrations in prepared films	1 mM, 2 mM, & 5 mM
8	Dye concentration in the solution (Chloroform)	0.0001 mol/L
9	Film Thickness	10 $\mu$ m
10	Single photon florescence Peak	Not applicable

### 3. NONLINEAR OPTICAL PROPERTIES OF DY-7

#### 3.1 Nonlinear Absorption Study (Open Aperture Z-Scan):

The block diagram of the experimental setup used for the open aperture Z-scan study of DY-7 doped in PMMA-MA film is shown in figure 3. A CW diode laser of wavelength 532nm (30 mW Green Light) is used as the light source. The Gaussian profiled laser beam is focused by a convex lens (L1) of focal length  $f$  ( $f = 3.5$  cm) to create a beam waist  $\omega_0$  of 15  $\mu$ m. The Rayleigh condition, diffraction length  $z_R = \pi\omega_0^2/\lambda > L$  is satisfied in this case so that the dye-doped sample is considered as a thin

medium, where  $L$  is the thickness of the sample and  $\lambda$  is the free space wavelength of the laser beam. The transmitted beam is collected by means of a convex lens and the output intensity measured using photo detector - digital power meter assembly.

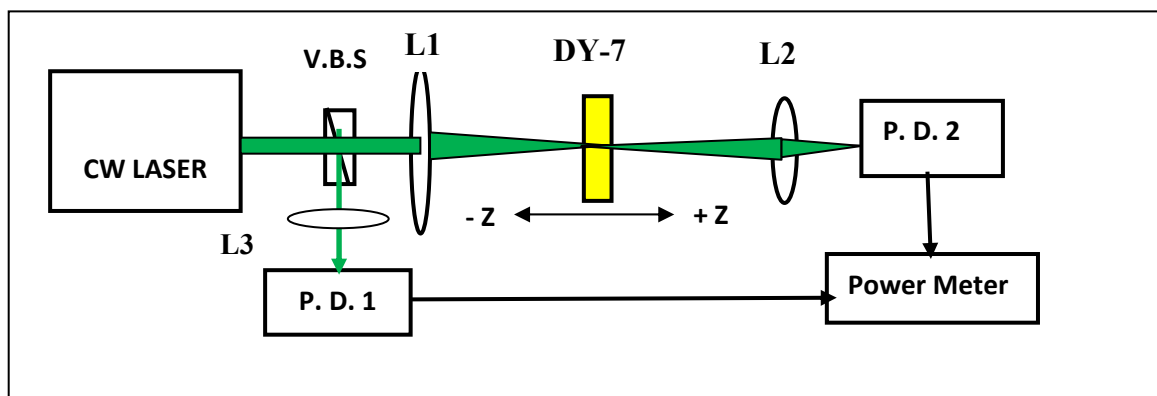


Figure 3: Experimental setup for open aperture Z scan for DY-7 samples.

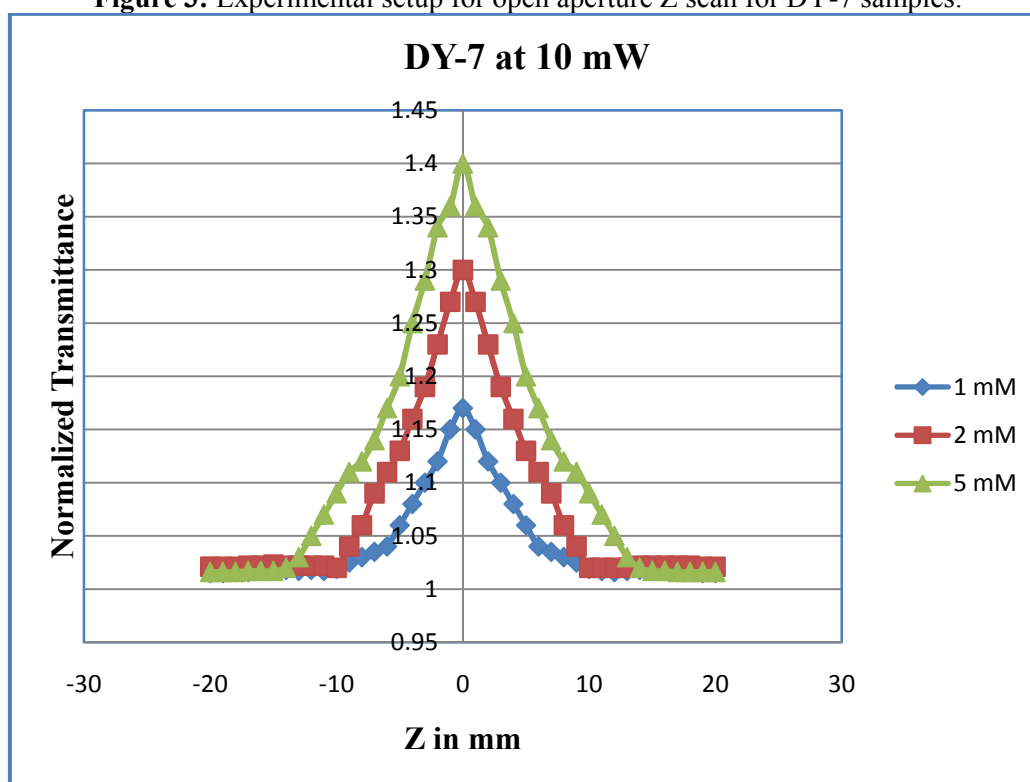


Figure 4 : CW Open aperture Z-scan plot of DY-7 at different dye concentrations using 532 nm, 10 mW laser beam.

The experimental setup used in the open aperture Z-scan is not sensitive to nonlinear refraction and hence can be used to determine the nonlinear absorption cross section of the materials. Such open aperture Z-scan trace is expected to be symmetric with respect to the focus where  $Z = 0$ , and at the focus, the minimum transmittance (e.g., multi-photon absorption/reverse saturation absorption) or a maximum transmittance (e.g., saturation of absorption) occurs. The nonlinear absorption coefficient ( $\beta$ ) can be estimated from Z-scan transmittance curve (Eq. 3). The diffraction length for the experimental setup is found to be 2.5 mm. The input power adjusted and noted by means of a convex lens (L3) and Photo detector (P.D.1). The DY-7 dye-doped polymer sample is translated across the focal region of lens L1 along the axial direction that is the direction of the propagation of the laser beam.

The Z- scan experiment is performed for DY-7 dye-doped PMMA-MA polymer films of the dye concentration 1mM, 2mM, and 5 mM using 532 nm laser beam at 10mW, 20 mW and 30 mW. The results are depicted in Fig. 4, to Fig. 6 respectively. In open aperture Z-scan, DY-7 has shown a decrease in transmittance with an increase in irradiance/input intensity due to reverse saturation absorption [16-17].

It is seen from the Z-scan plot that the DY-7 shows strong saturable absorption at low input intensity of laser beam. From the open aperture Z-scan, it is observed that due to nonlinear absorption, the transmittance of the DY-7 film is increased initially with an increase in intensity due to saturation of absorption and as input power increases, the saturation absorption (SA) is overtaken by reverse saturation absorption (RSA) as seen in figure 6. Such transformation from saturation absorption to reverse saturation absorption can be utilized as a principle for the construction of optical switches as well as optical limiters.

Based on open aperture Z-scan plots of DY-7 for different concentrations and at different input power, it is observed that :

(i) At low input power, saturation absorption (SA) increased with increase in the concentration of dye in the sample.

(ii) At the higher intensity of input light, DY-7 has shown reverse saturation absorption (RSA) so that saturation absorption (SA) of the sample is decreased.

From the graphs (figure 5 and figure 6), reverse saturable absorption (RSA) is seen in the open aperture Z-scan trace for DY-7 dye doped in PMMA-MA film as it shows minimum transmittance with increase in intensity of input laser beam. The nonlinear absorption coefficient  $\beta$  can be estimated from the open aperture Z-scan data, where

$$\beta = (2\sqrt{2} \Delta T) / (I_0 L_{eff}) \text{ ----- (3)}$$

Here,  $I_0$  is the intensity at the focal spot and is given by

$$I_0 = 2P_{peak} / \pi \omega_0^2 \text{ ----- (4)}$$

The effective length of the sample can be determined from the formula

$$L_{eff} = (1 - e^{-\alpha_0 L}) / \alpha_0 \text{ ----- (5)}$$

For low input intensity, the transmittance increases with the increase in excitation intensity and has a maximum value at the focus. As input power is increased, the sample has shown a decrease in transmittance which is the signature of reverse saturation absorption according to Sheik-Bahae's theory [18-19]. When reverse saturation absorption occurs, the absorption coefficient  $\beta$  is no longer a constant. Instead, it becomes a function of the excitation intensity as in the relation,

$$\alpha = \alpha_0 + I\beta \text{ ----- (6)}$$

here,  $\alpha$  is the total absorption coefficient of the material,  $\alpha_0$  is the linear absorption coefficient of the material,  $\beta$  is the nonlinear absorption coefficient of the material, and  $I$  is the input intensity of the laser beam.

By considering only third-order nonlinearities in the sample, the total refractive index of the sample ( $n$ ) becomes:

$$n = n_0 + n_2 I \text{ ----- (7)}$$

where  $\Delta n = n_2 I$  is change in refractive index,  $n_0$  is the linear refractive index of the sample,  $n_2$  is the nonlinear refractive index for third-order nonlinear sample, and  $I$  is the input intensity of the laser beam.

Reverse saturable absorption is seen in the open aperture Z-scan trace for DY-7 dye in PMMA-MA matrix as it shows minimum transmittance. The nonlinear absorption coefficient  $\beta$  can be estimated from the open aperture Z-scan data using eq. (3). The transmittance increases with the increasing input intensity and has a maximum value at the focus where  $Z=0$ , which is the sign of saturation absorption according to Sheik-Bahae's theory. In eq. (3),  $\Delta T$  is maximum transmittance at the focus (at  $Z = 0$ ). When saturation absorption occurs, the absorption coefficient  $\alpha$  will not become a constant. Instead, it will become a function of the excitation intensity as in the relation,  $\alpha = \alpha_0 + I\beta$  where  $\alpha_0$  is the linear absorption coefficient and  $\beta$  is its nonlinear counterpart. The experimental values of nonlinear absorption coefficient  $\beta$  at different dye concentrations for DY-7 dye-doped PMMA-MA films at 20 mW input power are listed in Table 3.

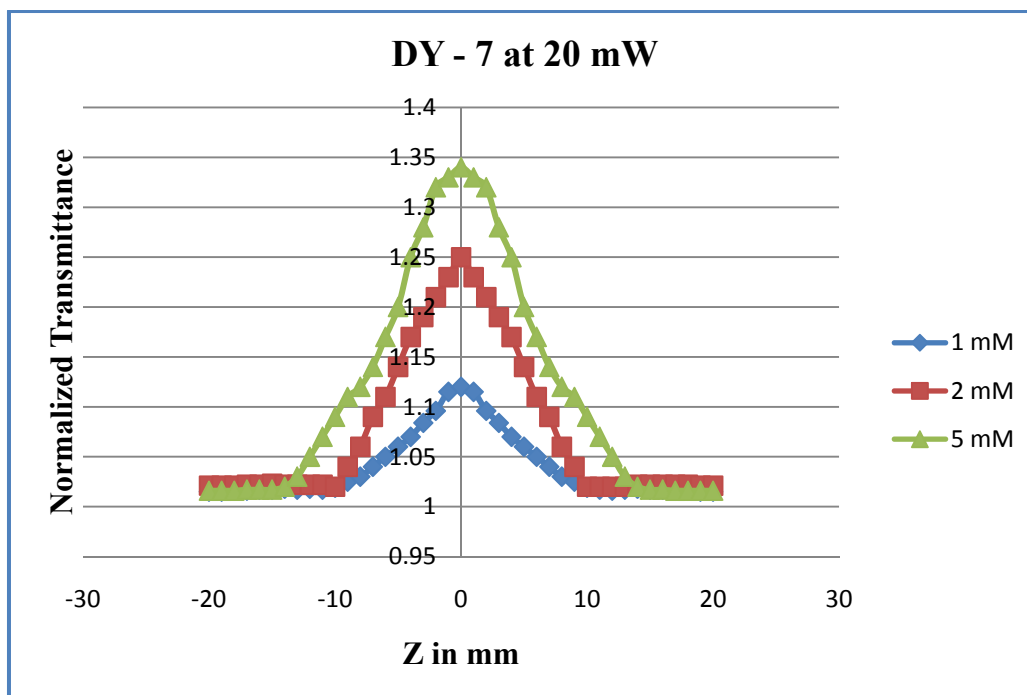


Figure 5: CW Open aperture Z-scan plot of DY-7 at different dye concentrations using 532 nm, 20 mW laser beam.

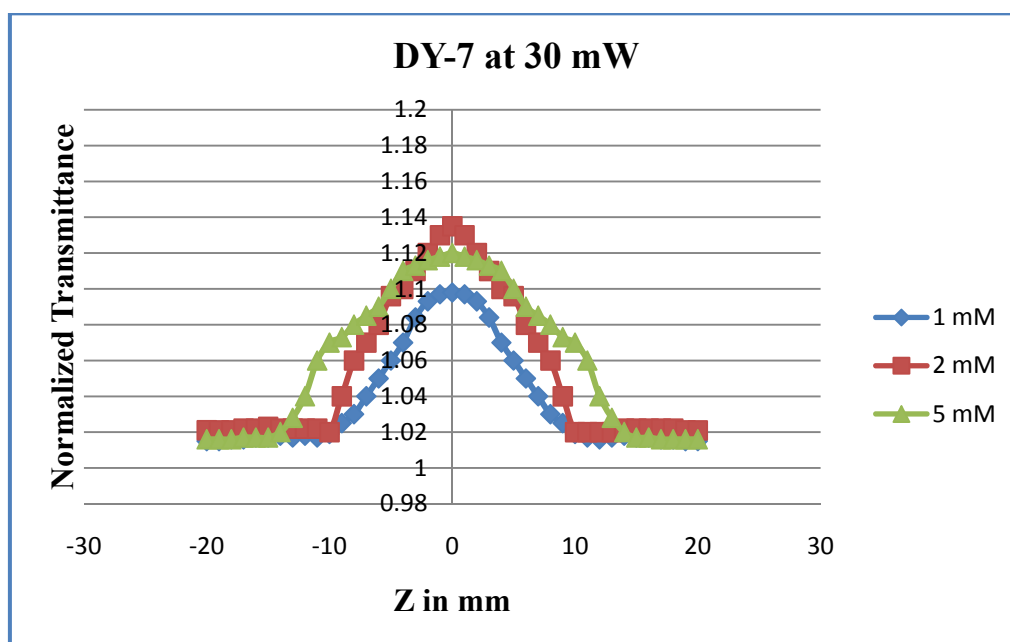


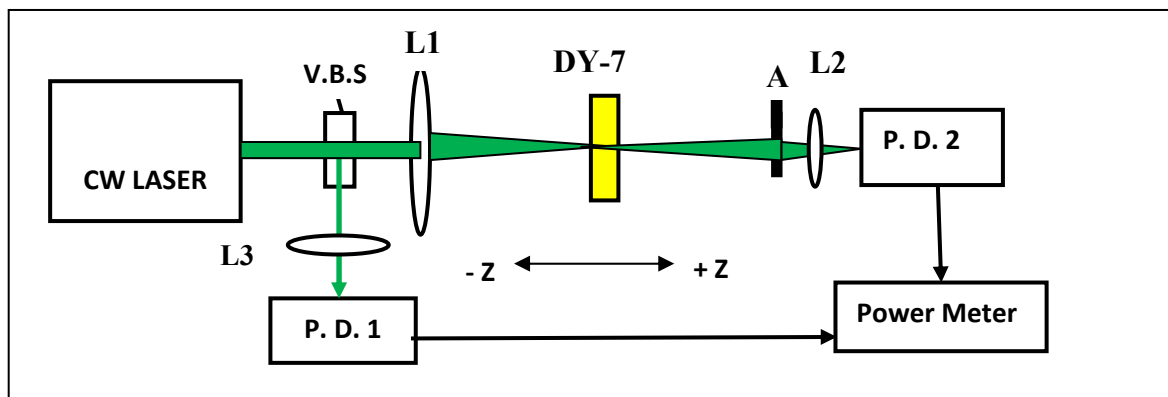
Figure 6 : CW Open aperture Z-scan plot of DY-7 at different dye concentrations using 532 nm, 30 mW laser beam.

Table 3 : Nonlinear absorption coefficient values for DY-7 dye-doped PMMA-MA film.

Wavelength (nm)	Concentration	$\Delta T$	$I_0$ (KW/cm <sup>2</sup> )	$\beta$ (cm/W) $\times 10^{-3}$
532	1 mM	0.12	3.5	- 0.97
532	2 mM	0.26	3.5	- 2.10
532	5 mM	0.34	3.5	- 2.75



### 3.2 Nonlinear Refraction Study (Closed Aperture Z-Scan) :



**Figure 7 :** Experimental setup of closed aperture Z scan for DY-7 sample.

The experimental set-up used for the closed aperture Z-scan technique is same as the set-up used for open aperture Z-scan except for the output beam from the DY-7 dye-doped sample is collected through an aperture of a fixed hole size instead of collecting entire output beam through collecting lens L2. The diode laser of wavelength 532 nm (30 mW Green Light) is used as the excitation source and the Gaussian beam is focused by means of a 3.5 cm focal length convex lens (L1), which produced a beam waist  $\omega_0$  of 15  $\mu\text{m}$ . The peak intensity of the incident laser beam is calculated as  $I_0 = 3.5 \text{ kW/cm}^2$  and the diffraction length ( $Z_R$ ) is calculated as 2.5 mm. The schematic of the experimental set up used is shown in figure 7. The dye sample is translated across the focal region along the axial direction that is the direction of the propagation laser beam. The intensity of transmitted light through an aperture kept in the far field is measured using photo detector and the digital power meter assembly. The closed aperture Z-scan plot between Z in mm and normalized transmittance for different dye concentration is shown in figure 8.

The observed peak amplitude followed by a valley amplitude of normalized transmittance curve in the closed aperture Z-scan plot shows that the sign of the refractive nonlinearity is negative, i.e. the DY-7 dye-doped sample is self-defocusing property. The self-defocusing effect is due to the local variation of the refractive index with temperature. The defocusing effect for the DY-7 dye in polymer film shown in figure 8 is attributed to a thermal nonlinearity originating from intensity dependent absorption of light at 532 nm. The change in transmittance  $\Delta T_{p-v}$  can be measured as the difference between the normalized peak transmittance value and valley transmittance value. Since the closed aperture transmittance is affected by the nonlinear refraction and absorption, the determination of nonlinear refraction is not easy and direct. To solve this, one has to separate the effect of nonlinear absorption from that of the nonlinear refraction. The simple solution to get purely refractive nonlinearity is to divide the closed aperture transmittance data by the corresponding open aperture transmittance data [20-21].

The pure nonlinear refractive Z-scan curves are shown in figure 9 for the Dy-7 dye in the polymer film. In order to know the contribution from pure PMMA-MA polymer film to the observed nonlinear response, the Z-scan is performed on pure film without DY-7 dye doping. Neither nonlinear absorption nor nonlinear refraction is observed. The nonlinear refractive index  $n_2$  can be calculated using the Eq. (8). The experimental values of nonlinear phase shift and nonlinear refractive index for DY-7 dye-doped PMMA-MA film are listed in Table 4.

The nonlinear refractive index  $n_2$  can be calculated using the formula  $n_2 = \frac{\Delta\phi \lambda}{2\pi I_0 L_{eff}}$  ----- (8)

and  $|\Delta\phi| = \Delta T_{(p-v)} / [0.406 (1-S)^{0.25}]$  ----- (9)

where  $\Delta T_{p-v}$  is the peak-valley transmittance difference from the closed aperture plot,  $|\Delta\phi_0|$  is the on axis nonlinear phase-shift and S is the aperture linear transmittance given by  $S = [1 - \exp(-2r_a^2/w_a^2)]$  where  $r_a$  is the aperture radius and  $w_a$  is the beam radius at the aperture.  $S=1$  for open aperture configuration and S is adjusted to 0.5 for our closed aperture configurations. In Eq. (8),  $I_0$  is the intensity at the focal spot as per Eq. (4) and  $L_{eff}$  is the effective length of the sample and is given by



Eq. (5).

The change in refractive index  $\Delta n$  can be calculated using the formula,

$$\Delta n = n_2 I_0 \quad \text{----- (10).}$$

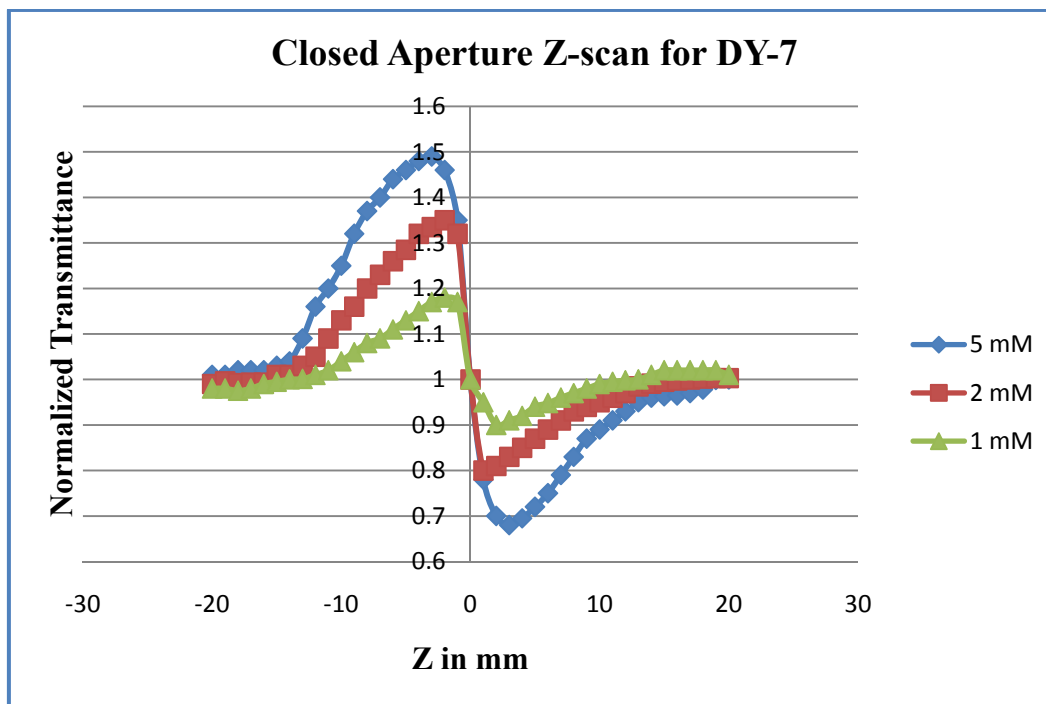


Figure 8: Nonlinear refraction of DY-7 using CW 532 nm 20 mW laser beam.

Table 4: Nonlinear refractive index and nonlinear phase shift for DY-7 film.

Wavelength (nm)	Dye Concentration	$\Delta T_{p-v}$	$I_0 \text{ mW/cm}^2$	$\Delta\phi_{rad}$	$n_2 (\text{cm}^2/\text{mW}) \times 10^{-8}$
532	1 mM	0.29	3.5	0.8494	- 2.48
532	2 mM	0.55	3.5	1.6110	- 4.71
532	5 mM	0.82	3.5	2.4010	-7.02

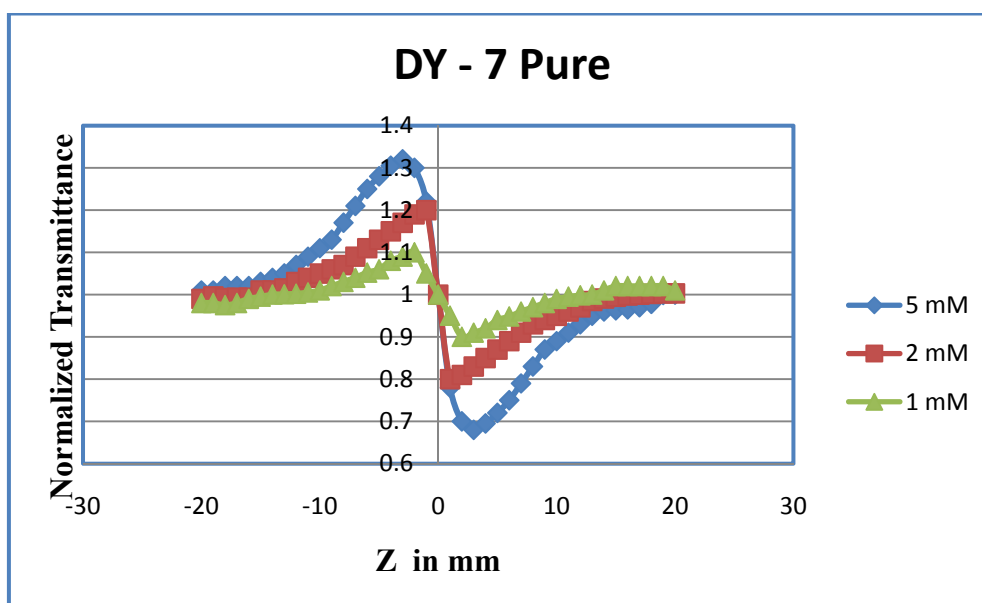


Figure 9: Pure nonlinear refraction plot of DY-7 using CW 532 nm 20 mW laser beam.

### 3.3 Third-order Nonlinear Optical Susceptibility $|\chi^{(3)}|$ :

The Z- scan plot of DY-7 dye in PMMA-MA polymer film, show a pre-focal transmittance maximum (peak) followed by a post-focal transmittance minimum (valley). This indicates that DY-7 has a negative nonlinearity due to self-defocusing. Self-defocusing is due to variation in refractive index with the thermally agitated dye molecules along with reverse saturation absorption. The nonlinear refractive index  $n_2$  can be calculated using the equation (8) and change in refractive index,  $\Delta n$  can be calculated using equation (10). Experimentally determined nonlinear refractive index  $n_2$  can be used to find the real part and imaginary part of the third-order nonlinear optical susceptibility  $[\chi^3]$  using Eqs. (11), (12), and (13).

Experimentally determined nonlinear refractive index  $n_2$  can be used to find the real part of the third-order nonlinear optical susceptibility  $[\chi^3]$  according to the following relation [22-23].

$$\text{Re } \chi^{(3)} = |\chi^3| = \frac{10^{-4} \epsilon_0 (n_0^2) c^2 n_2}{\pi} \quad (\text{cm}^2/\text{W}) \quad \text{-----} \quad (11)$$

Experimentally determined nonlinear absorption coefficient  $\beta$  can be used to find the imaginary part of the third-order nonlinear optical susceptibility  $[\chi^3]$  according to the following relation

$$\text{Im } \chi^{(3)} = \frac{10^{-2} \epsilon_0 (n_0^2) c^2 \lambda}{4\pi^2} \quad (\text{cm}^2/\text{W}) \quad \text{-----} \quad (12)$$

The absolute value of the third-order nonlinear optical susceptibility is given by the relation

$$|\chi^3| = [(\text{Re } \chi^{(3)})^2 + (\text{Im } \chi^{(3)})^2]^{1/2} \quad \text{-----} \quad (13)$$

Where  $\epsilon_0$  is the vacuum permittivity and  $C$  is the light velocity in vacuum.

The nonlinear parameters, such as nonlinear refractive index ( $n_2$ ), change in refractive index ( $\Delta n$ ), the nonlinear absorption coefficient ( $\alpha$ ) and nonlinear susceptibility ( $\chi^{(3)}$ ) are calculated and listed in Table 5.

Table 5: Third harmonic susceptibility for DY-7 film.

Wavelength (nm)	Concentration	$n_2$ (cm <sup>2</sup> /W) $\times 10^{-8}$	$\Delta n$ ( $\times 10^{-4}$ )	$[\chi^3]$ (esu) $\times 10^{-6}$
532	1 mM	- 2.48	-0.868	5.12
	2 mM	-4.71	-1.649	6.82
	5 mM	-7.02	-2.457	8.52

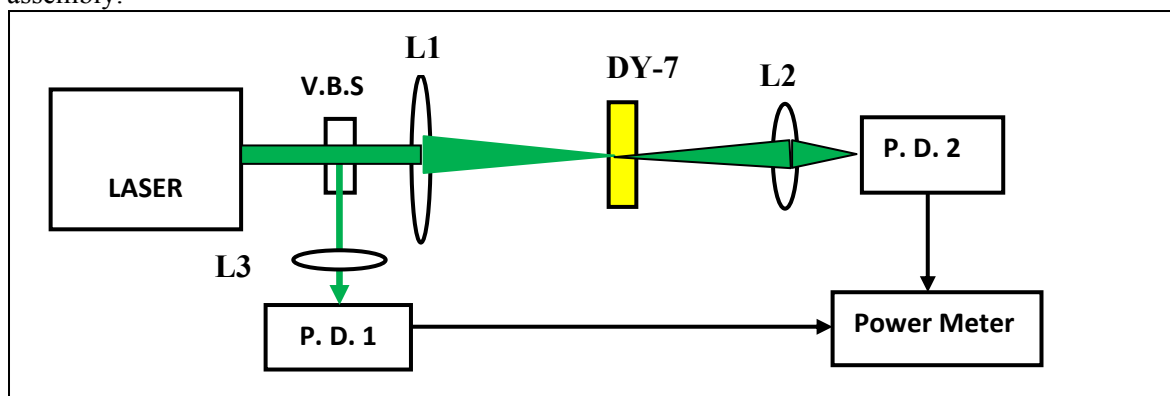
## 4. CW OPTICAL LIMITING STUDY

The optical limiting effect of the DY-7 dye-doped polymer film is analysed by means of using a 532 nm, 30 mW CW semiconductor diode laser beam (30 mW Green Light). Two experimental setups are used for the demonstration of optical limiting. In the first experimental setup, the dye sample is placed in the focus of the focusing lens L1 of Z-scan setup (figure 10). The emergent beam from the dye sample is collected to a photo detector by means of a collecting lens L2 to measure the output power. By fixing the sample position at the focus, the input power is varied and output power is noted. Such experimental setup is named as *Optical limiting without an aperture* or **Type 1 optical limiting**. This type of optical limiting study will take care of nonlinear absorption property of the dye sample. In the second experimental setup, an aperture of fixed hole size is used between the dye sample and the collecting lens & photo detector. The dye sample film is kept at a position where the transmitted intensity shows a valley in closed aperture Z-scan curve [24-26]. The input laser beam intensity is varied and the corresponding output laser beam intensity values are noted by a photo detector and power meter assembly. Such experimental setup is named as *Optical limiting with an aperture* or **Type 2 optical limiting**. This type of optical limiting study will take care of nonlinear refraction property of the dye sample.

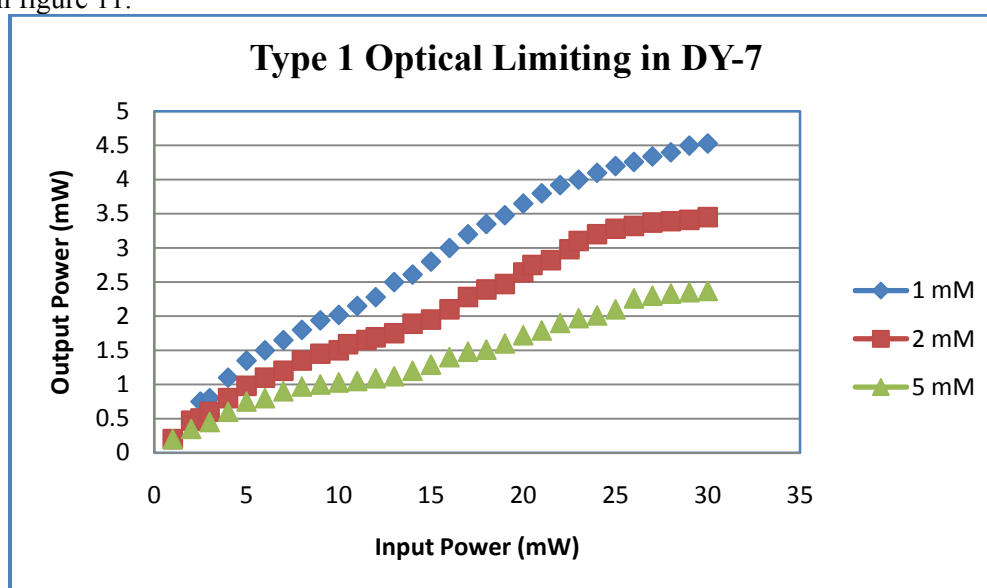
### Case (1): Optical Limiting without Aperture (Type 1):

The pure nonlinear absorption property of the dye sample is measured using this method of optical limiting without aperture at the output side (Type-1 optical limiting). The entire light beam transmitted through the sample is focused by a collecting lens to the photo detector-power meter assembly. The optical limiting effects of the DY-7 dye-doped PMMA-MA films are studied by using a CW laser source. The experimental set-up for the demonstration of type-1 optical limiting is shown in figure 10. The dye sample is kept fixed at the focal point of the convex lens L1 of open aperture Z-

scan setup. The input power is changed by means of a variable beam splitter (VBS). By means of a convex lens, the output light beam intensity is measured using a photo detector-power meter assembly.



**Figure 10:** Experimental setup for Optical limiting (Type 1) due to pure absorptive nonlinearity. The input light intensity is increased systematically and the corresponding output intensity is measured by a photo detector. The output power is measured using a power meter. The experiment is performed at different input power and the corresponding output power of transmitted beam is noted and a graph is drawn between input power and output power for different dye concentrations and is shown in figure 11.



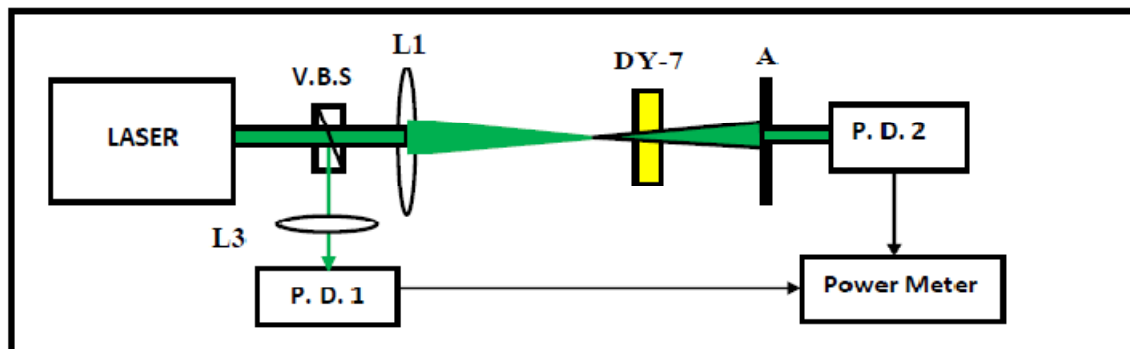
**Figure 11 :** Open aperture (Type 1) Optical limiting behavior of DY-7 dye-doped PMMA-MA film at CW 532 nm.

In this case, the intensity of the transmitted beam at the output is observed to change linearly at low incident input intensity but found to saturate at higher value of input intensities due to nonlinear reverse saturation absorption and thermal lensing effect. Hence, after a certain threshold value of the input light intensity, the nonlinear absorption of the DY-7 dye sample becomes dominant, resulting in a limiting of the intensity of output beam. Thus the transmittance recorded by the photo detector remained almost constant showing a saturation region.

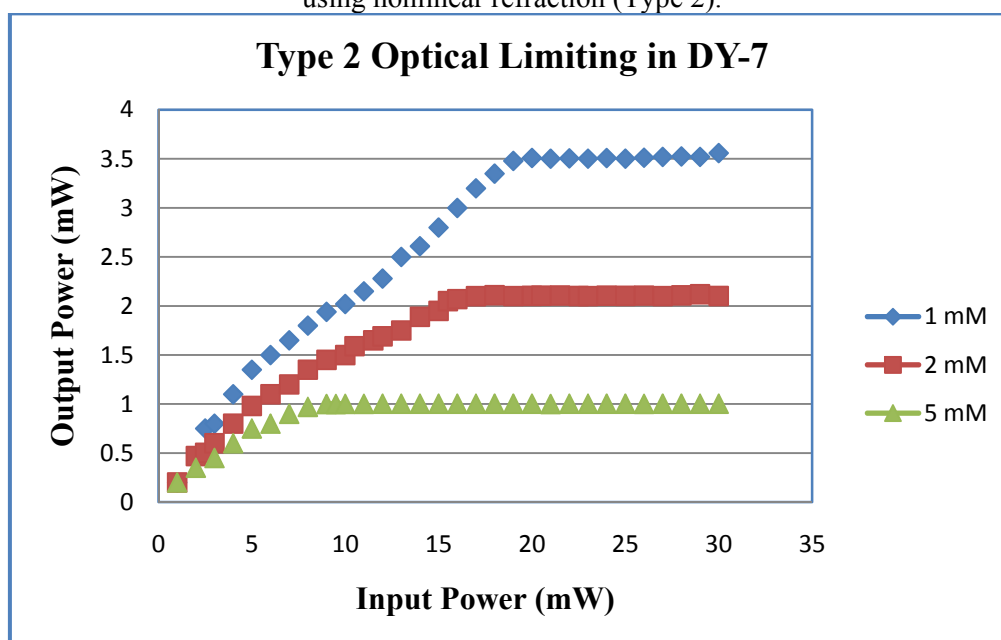
#### Case (2) Optical Limiting with Aperture (Type 2) :

The pure nonlinear refraction property of the dye sample is measured using this method of optical limiting with an aperture at the output side (Type-2 optical limiting). The light beam transmitted through the sample is passed through an aperture A of fixed diameter and then passed through a collecting lens L2 to the photo detector-power meter assembly. The optical limiting effects of the DY-7 dye-doped PMMA-MA films are studied by using a CW laser source. The experimental set-up for

the demonstration of type-2 optical limiting is shown in figure 12. The dye sample is kept at the position where the transmitted intensity shows a valley in the closed aperture Z-scan curve. The experiment is performed at different input power and the corresponding output power of transmitted beam is noted and a graph is drawn between input power and output power for different dye concentrations and is shown in figure 13.



**Figure 12:** Experimental setup for closed aperture Optical limiting for DY-7 doped PMMA-MA using nonlinear refraction (Type 2).



**Figure 13 :** Closed aperture (Type 2) Optical limiting behavior of DY-7 dye-doped PMMA-MA film at CW 532 nm.

In this case of DY-7 doped PMMA-MA film with defocusing nonlinearity (negative nonlinearity), the intensity of transmitted output laser beam is observed to change linearly with low values of input intensity, but found to saturate after a certain threshold value of input power due to the reason that the dye sample starts defocusing the transmitting beam, resulting blocking of the part of the beam by the aperture A placed in between the dye sample and the collecting lens before the photo-detector. Thus the light intensity received by the photo detector remained almost constant showing a saturation region as shown in figure 13.

## 5. RESULTS & DISCUSSION

In this section, we have studied the nonlinear optical properties of organic dye material DY-7 in PMMA-MA polymer matrix for their nonlinear absorption using open aperture Z-scan, nonlinear refraction using closed aperture Z-scan, and for optical limiting using type 1 and type 2 optical limiting configurations. Based on our study we observed that the nature of nonlinearity shown by DY-7 dyes depends on the intensity of input laser light beam. At the lower intensity of input beam, the

saturation absorption became prominent and with an increase in input intensity further, the excited state absorption in the form of reverse saturation absorption or thermally induced absorption became prominent.

It is also observed from other studies [27-37] that the value of  $\Delta T_{p-v}$  has increased in the case of dye-doped polymer films when compared to the dyes in the solvent. This shows that the change in refractive index due to reverse saturation absorption as well as thermal effect in solid media is larger compared to that of the liquid media. The heat dissipation rate in the liquid is more than that in the film because of the poor thermal conductivity of the polymer film. This leads to increase in temperature in solid media compared to temperature rise in liquid media which might increase nonlinear absorption. So, the nonlinear refractive index change is more in solid media than in liquid media.

In this study, it is observed that the value of the nonlinear refractive index and nonlinear absorption coefficient of DY-7 dye-doped polymer films depends on the concentration of the dyes in the polymer film. It is seen that there is an increasing trend in values of  $n_2$ ,  $\beta$ , and  $|\chi^{(3)}|$  as the concentration increases. This may be explained by the fact that the number of dye molecules increases when the concentration increases, more dye molecules get involved in the nonlinear process (two-photon absorption/reverse saturation) due to thermal agitation resulting in an enhanced effect. Thus Z-scan measurements indicate that these dyes exhibit large nonlinear optical properties.

**Table 6 :** Nature of Nonlinearity of DY-7 dye doped in PMMA-MA polymer films.

S. No.	Dye	Medium	Nature of Nonlinearity
1	Disperse Yellow DY-7	PMMA-MA polymer film	Negative Nonlinearity Defocusing effect

**Table 7 :** Nonlinear parameters for DO-7 dye-doped films at 532 nm.

S. No.	Parameter	Dye concentration	DY-7
1	$\beta$ ( $\times 10^{-3}$ ) (cm/W)	1 mM	-0.97
		2 mM	-2.10
		5 mM	-2.75
2	$n_2$ ( $\times 10^{-7}$ ) (cm <sup>2</sup> /W)	1 mM	-0.248
		2 mM	-0.471
		5 mM	-0.702
3	$\Delta n = n_2 I_0$ ( $\times 10^{-4}$ )	1 mM	-0.868
		2 mM	-1.649
		5 mM	-2.457
4	$ \chi^3 $ ( $\times 10^{-6}$ ) (e.s.u.)	1 mM	5.12
		2 mM	6.82
		5 mM	8.52

The mechanism responsible for type 1 optical limiting is mainly attributed to reverse saturation absorption in case of DY-7, which further increased with thermally induced nonlinearity. The defocusing effect observed in DY-7 dye-doped samples under CW illumination is utilized to demonstrate type 2 optical limiting action. Based on their nonlinear refractive index, these dyes in PMMA-MA matrix behave as good optical limiters even at low powers. These results are quite impressive and encouraging for possible applications in nonlinear optical devices.

Table 8 lists details on various optical limiting regions in DY-7 dye-doped Polymer films at 532 nm CW laser beam. It is found that both type 1 and type 2 optical limiting effects show an increase in limiting action with increasing the concentration of the DY-7 dye in the polymer film. The optical limiting responses of the low dye concentration films are generally much weaker than those of high dye concentrated films. This shows that the number density of dye molecules in the polymer matrix along the path of the laser beam is the deciding factor to fix output clamping level.

**Table 8 :** Optical limiting Regions in Dye-doped Polymer films at 532 nm CW laser beam.

S. No.	Sample	Dye Concentration	Linear Region (mW)	Active Region (mW)	Saturation Region (mW)
1	DY-7 in PMMA-MA Type 1	1 mM	1 – 5 mW	5 – 23 mW	23 mW onwards
		2 mM	1 – 8 mW	8 – 25 mW	25 mW onwards
		5 mM	1 – 10 mW	10 – 26 mW	26 mW onwards
2	DY-7 in PMMA-MA Type 2	1 mM	1 - 18 mW	18 - 19 mW	19 mW onwards
		2 mM	1 – 15 mW	15 – 17 mW	17 mW onwards
		3 mM	1 – 8 mW	8 – 9 mW	9 mW onwards

From Table 9, it can be seen that the optical power limiting threshold is inversely proportional to the dye concentration in the film. The limiting experiment shows that as the concentration increases, a reduction in linear transmittance as well as the output clamping level. The experimentally determined optical limiting saturated output power values at different dye concentrations are shown in Table 10. The results are comparable to some of the reports of low power optical limiting [38-54].

In the case of type 2 optical limiter with aperture, as observed in our experiment and in other published results, it is seen that at the valley positions, the limiter works at low input powers as the self-focusing/self-defocusing effect is increased by the thermal effect due to the absorptive properties of the dye used in polymer matrix. Thus it can be suggested that the best position for a dye sample, when used for optical limiting based on Type 2 self-focusing/self-defocusing position is at the valley point of the Z-scan curve.

**Table 9 :** Concentration dependence of limiting Input threshold of DY-7 dye-doped in PMMA-MA films.

S. No.	Sample	Dye Concentration (mM)	Type 1 Optical Limiting Threshold (mW)	Type 2 Optical Limiting Threshold (mW)
1	DY-7 in PMMA-MA	1 mM	23	19
		2 mM	25	17
		5 mM	26	09

**Table 10 :** Concentration dependence of saturated output power in DY-7 dye-doped PMMA-MA films.

S. No.	Sample	Dye Concentration (mM)	Type 1 Optical Limiting Saturated Output Power (mW)	Type 2 Optical Limiting Saturated Output Power (mW)
1	DY-7 in PMMA-MA	1 mM	4.3	3.50
		2 mM	3.5	2.10
		5 mM	2.2	0.95

## 6. CONCLUSION

The nonlinear absorption, nonlinear refraction properties of DY-7 dye-doped in Polymethyl methacrylate methacrylic acid (PMMA-MA) polymer film are studied at low power CW laser beam of 532 nm using the Z-scan experimental method. The optical limiting properties of these films are also studied by increasing input power at different dye doping concentrations. It is found that the type of nonlinear absorption depends on the intensity of input beam. DY-7 has shown saturation absorption at lower input irradiance and reverse saturation absorption at higher irradiance. Optical limiting studies using type 1 and type 2 setups is carried out and is found that type 2 has shown better limiting characteristics for DY-7 doped PMMA-MA polymer films. Hence DY-7 dye-doped in PMMA-MA polymer film can be a potential candidate for the third harmonic property based photonic devices.



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